

# Emergence

Symmetry Breaking,  
Topological Defects  
and Effective Theories

Lecture I

9 AM

27 June 2011



# A General Definition of Emergence

A system has emergent properties when an effective theory of the system at some scale, or level of organization, is qualitatively different from the lower-level theory

# A General Definition of Emergence

A system has emergent properties when an **effective theory** of the system at some **scale**, or **level of organization**, is **qualitatively different** from the lower-level theory

# One Kind of Emergence

A system is emergent when the **symmetries** of the lower-level theory are **violated** under aggregation

# Topics

## Symmetry

Permutations, Shifts, and Finite Group Theory  
Invariances of Equations of Motion  
Continuous Symmetries  
Semigroups & “approximate” Symmetries

## Symmetry Breaking

Essential vs. Spontaneous  
Navier-Stokes & Turbulent Symmetry Breaking  
Symmetry Restored?

## Phase Transitions

Ising and XY Models  
Annealing vs. Domain Wall Formation  
Effective Theories for Defects

## Emergence Defined

## Signatures of Emergence in Animal Society

# Essential Symmetry Breaking

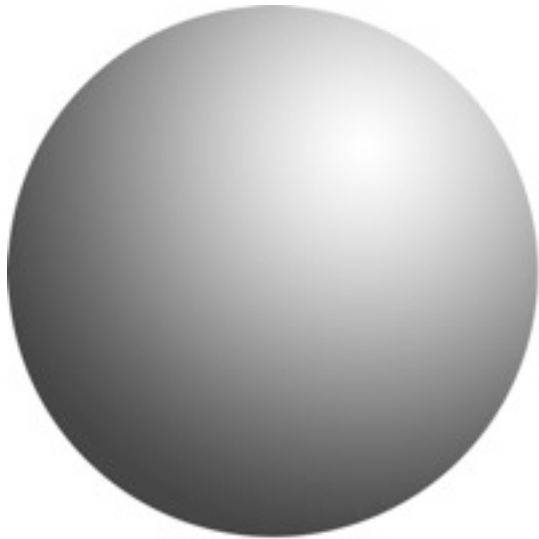
$$H(\sigma_i, \sigma_j) \sim \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$



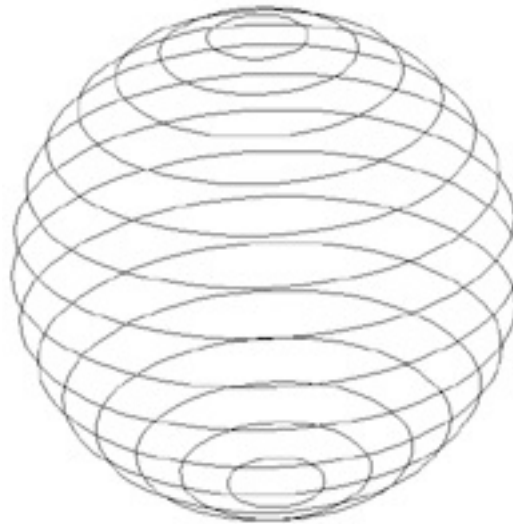
$$H(\sigma_i, \sigma_j) \sim \begin{pmatrix} 2 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$S_3 \rightarrow Z_2$$

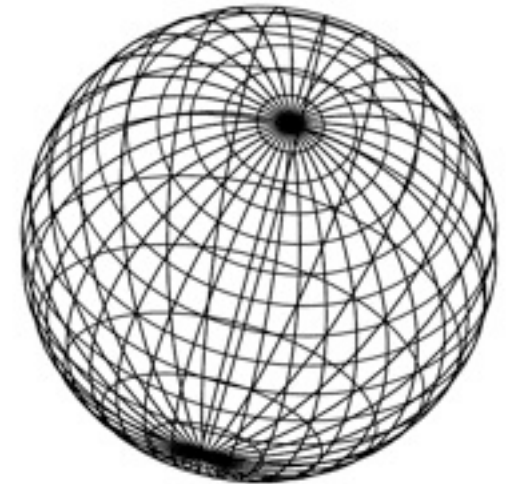
# Continuous Symmetries



all rotations in 3 dimensions  
a.k.a.,  $O_3$



only rotations in 2 dimensions  
a.k.a.,  $O_2$



*discrete* rotations in 2  
dimensions  
a.k.a.,  $Z_n$   
(where  $n$  is the number of  
longitude marks.)

$$O_3 \rightarrow O_2 \rightarrow Z_2$$

# Traffic Jam without Bottleneck

Experimental evidence  
for the physical mechanism of forming a jam

Yuki Sugiyama, Minoru Fukui, Macoto Kikuchi,  
Katsuya Hasebe, Akihiro Nakayama, Katsuhiro Nishinari,  
Shin-ichi Tadaki and Satoshi Yukawa

## Movie 1

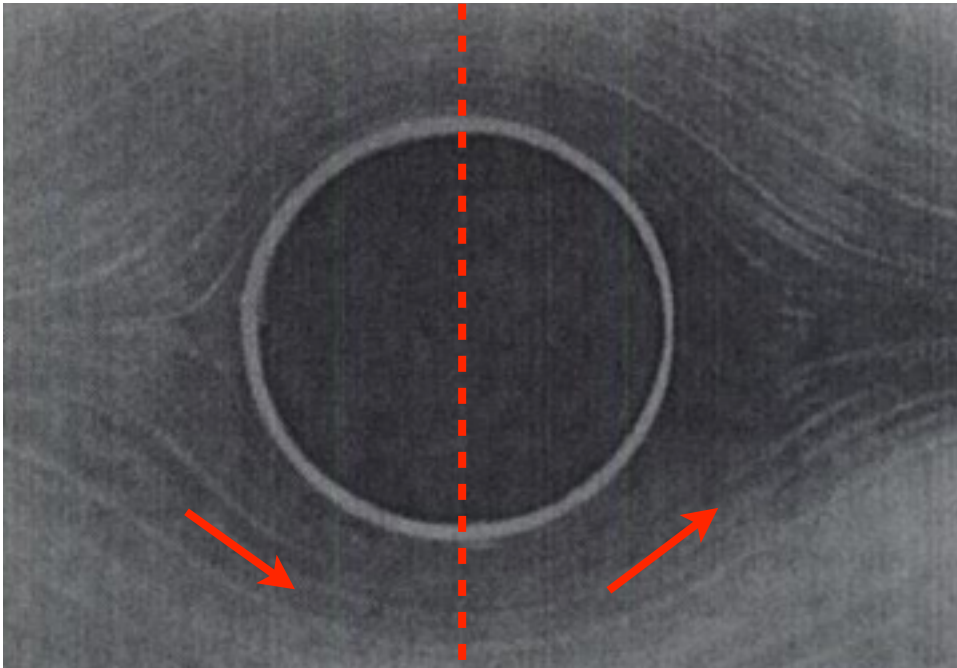


The Mathematical Society of Traffic Flow

# Navier-Stokes Equation

$$\frac{\partial v_i}{\partial t} \quad \text{[Grey Box]} = -\frac{\partial p}{\partial x_i} + \nu \sum_j \frac{\partial^2}{\partial x_j^2} v_i$$

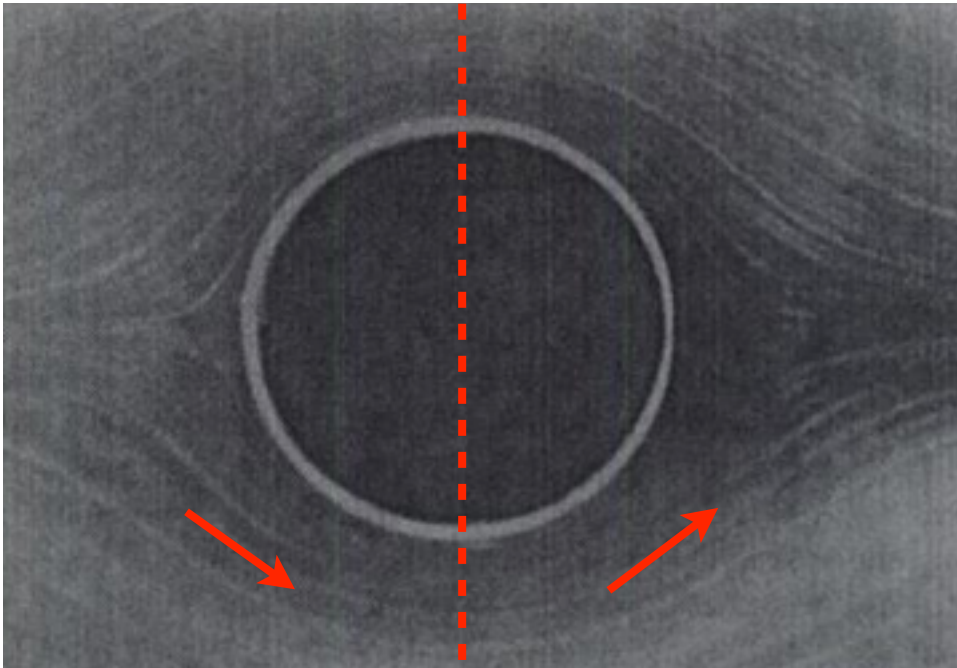
Symmetry: flip x position, y velocity



# Navier-Stokes Equation

$$\frac{\partial v_i}{\partial t} + \sum_j v_j \frac{\partial}{\partial x_j} v_i = -\frac{\partial p}{\partial x_i} + \nu \sum_j \frac{\partial^2}{\partial x_j^2} v_i$$

Symmetry: flip x position, y velocity

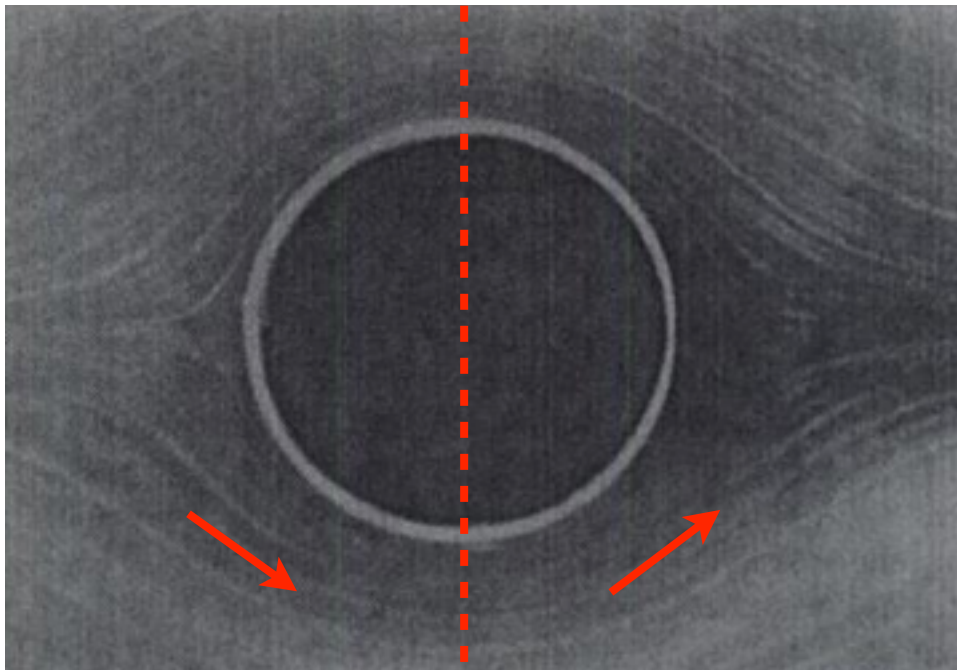


Uriel Frisch, *Turbulence*

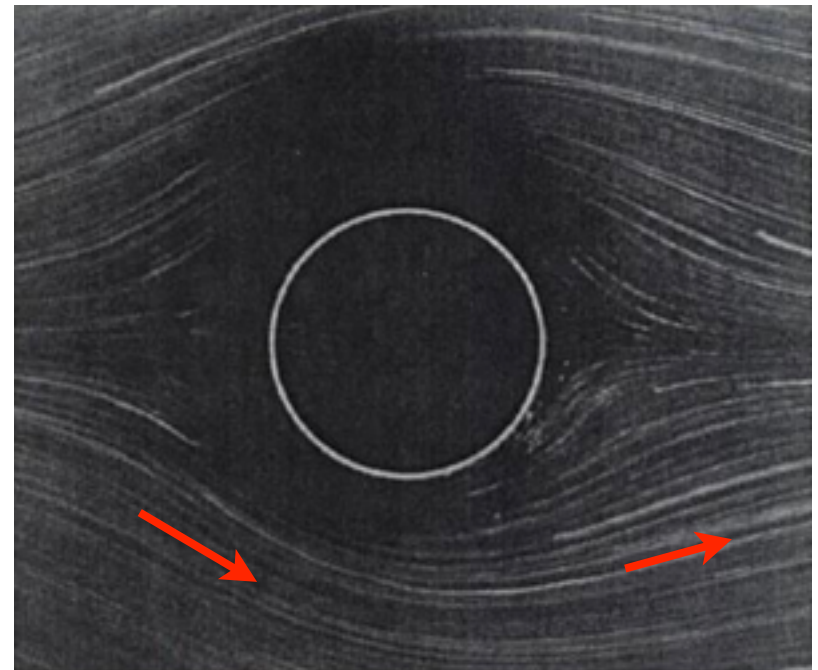
# Navier-Stokes Equation

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Symmetry: flip x position, y velocity



...upon turning on non-linear term

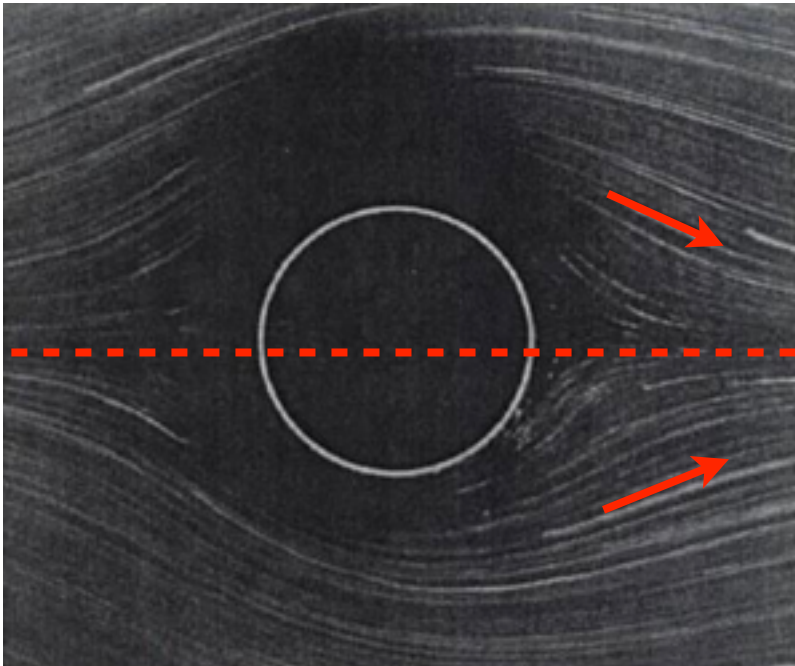


Uriel Frisch, *Turbulence*

# Navier-Stokes Equation

$$\frac{\partial v_i}{\partial t} + \sum_j v_j \frac{\partial}{\partial x_j} v_i = -\frac{\partial p}{\partial x_i} + \nu \sum_j \frac{\partial^2}{\partial x_j^2} v_i$$

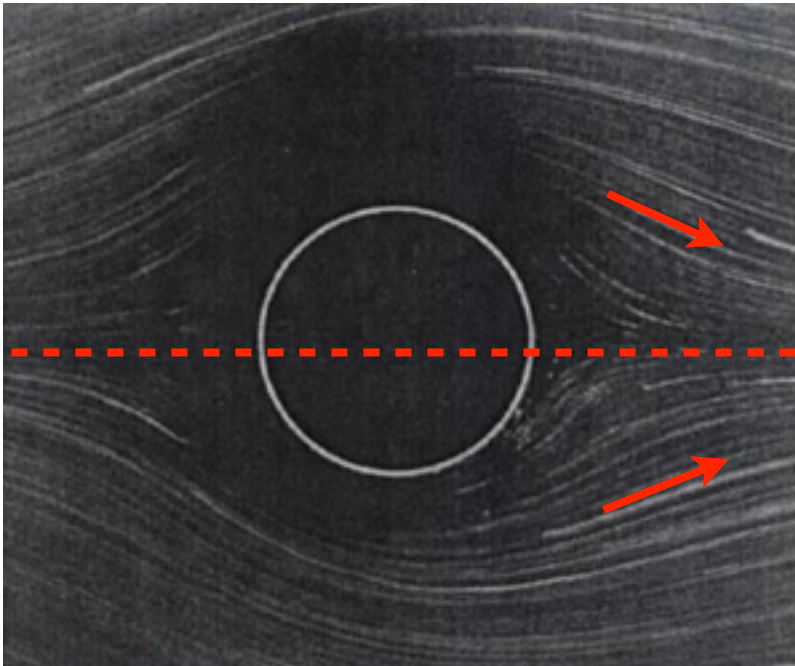
y symmetry still exists...



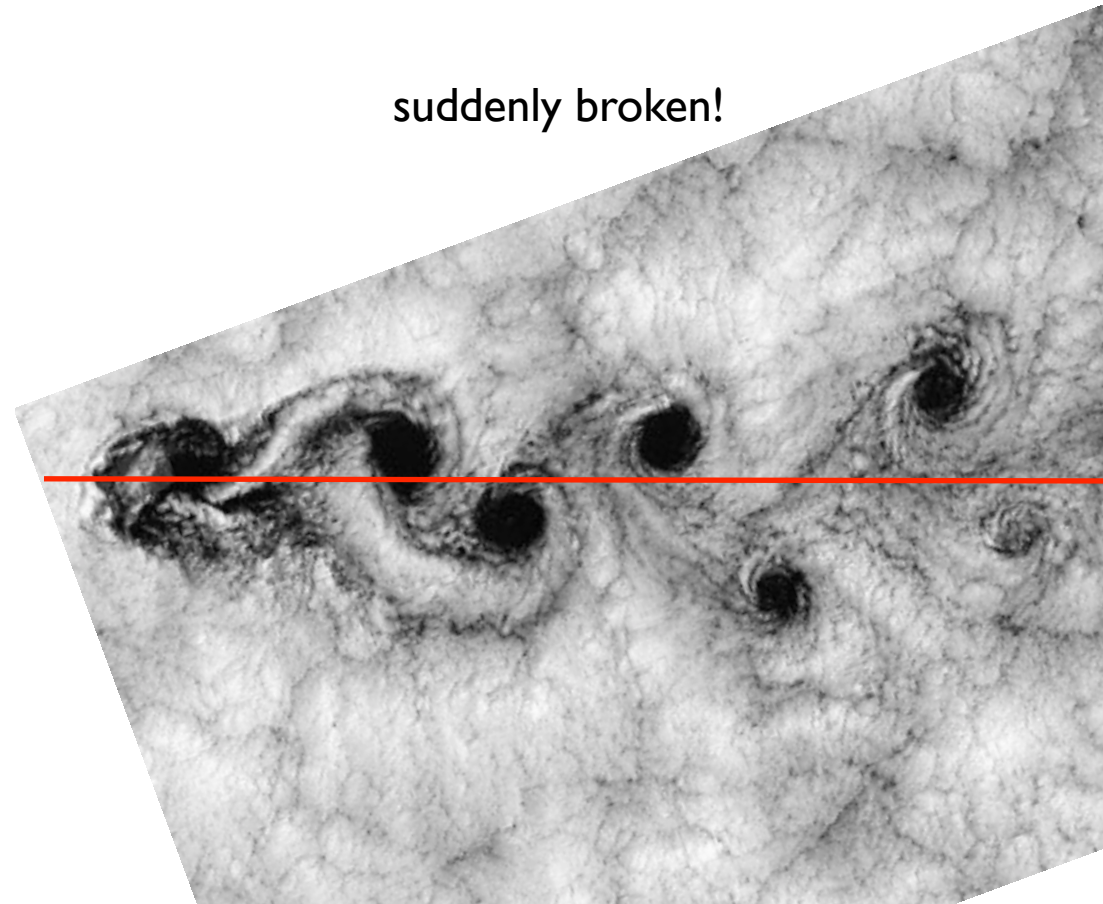
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y symmetry still exists...



suddenly broken!



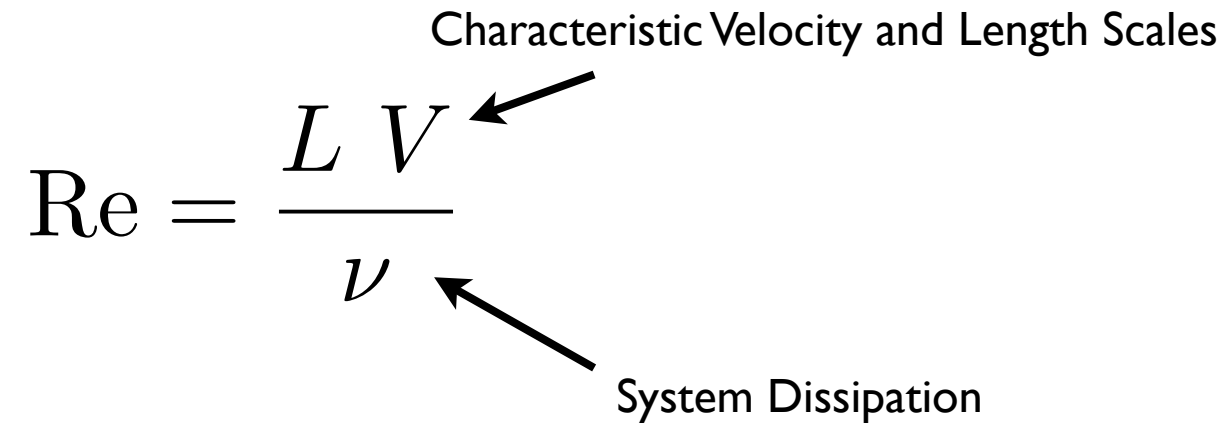
# Control Parameter

Reynold's Number

$$\text{Re} = \frac{L V}{\nu}$$

Characteristic Velocity and Length Scales

System Dissipation



Pablo Navarrete Michelini, University of Chile

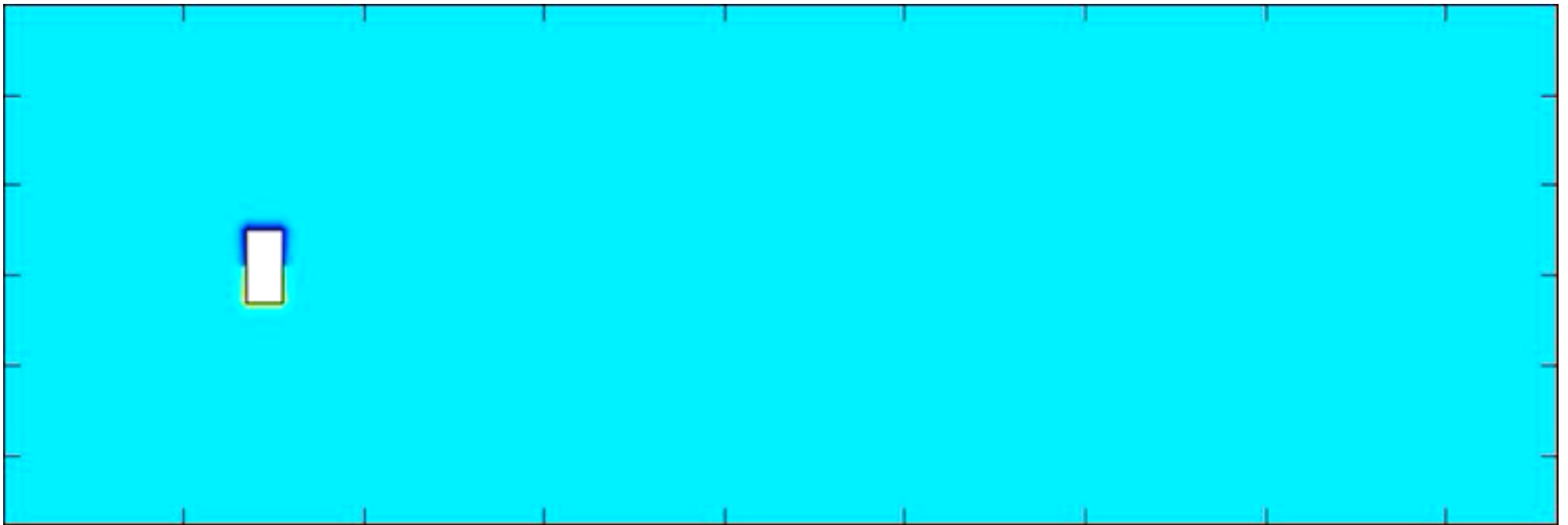
# Control Parameter

Reynold's Number

$$\text{Re} = \frac{L V}{\nu}$$

Characteristic Velocity and Length Scales

System Dissipation



Pablo Navarrete Michelini, University of Chile

# What is your symmetry?

discrete? Fully symmetric? Cyclic? A pair of cycles? A set of Nash Equilibria?  
continuous? A real number? A vector?  
[a fundamental mechanism question]

# What is your interaction?

pair wise? (neighbours on a graph, cars in a line?) Mean field? (Well mixed population?)  
Higher-order? (Cooperative effects, spatial structures...)  
[a fundamental mechanism question]

# What is your control parameter?

Noise? Dissipation? Penalty for failure?  
Strength of interaction? Of pairwise:triplet interaction ratio?  
[when the system shows emergence]

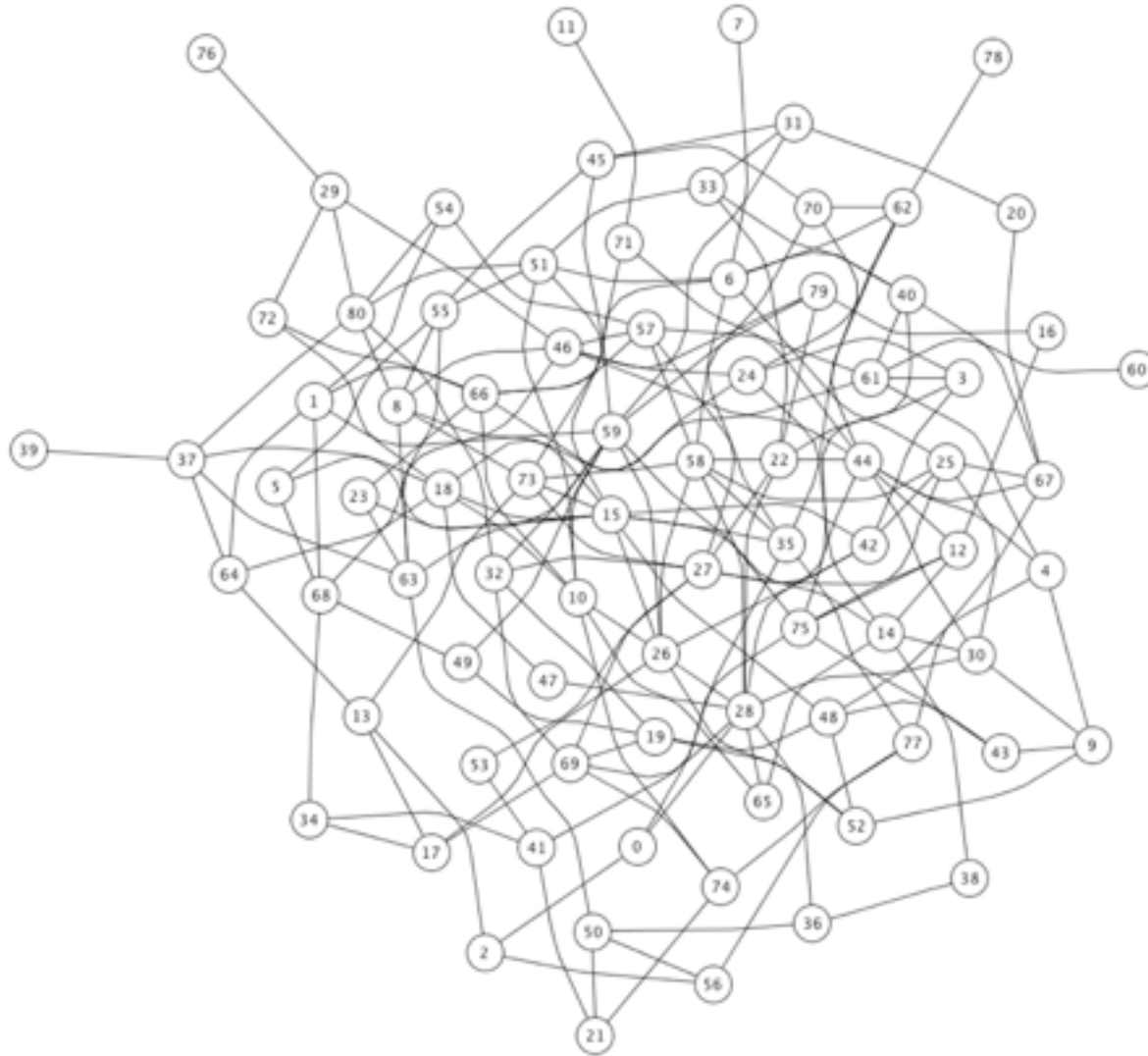


# The Cost Function of the Ising Model

$$H[\{\sigma_i\}] = - \sum_{i,j} J_{ij} \sigma_i \sigma_j$$

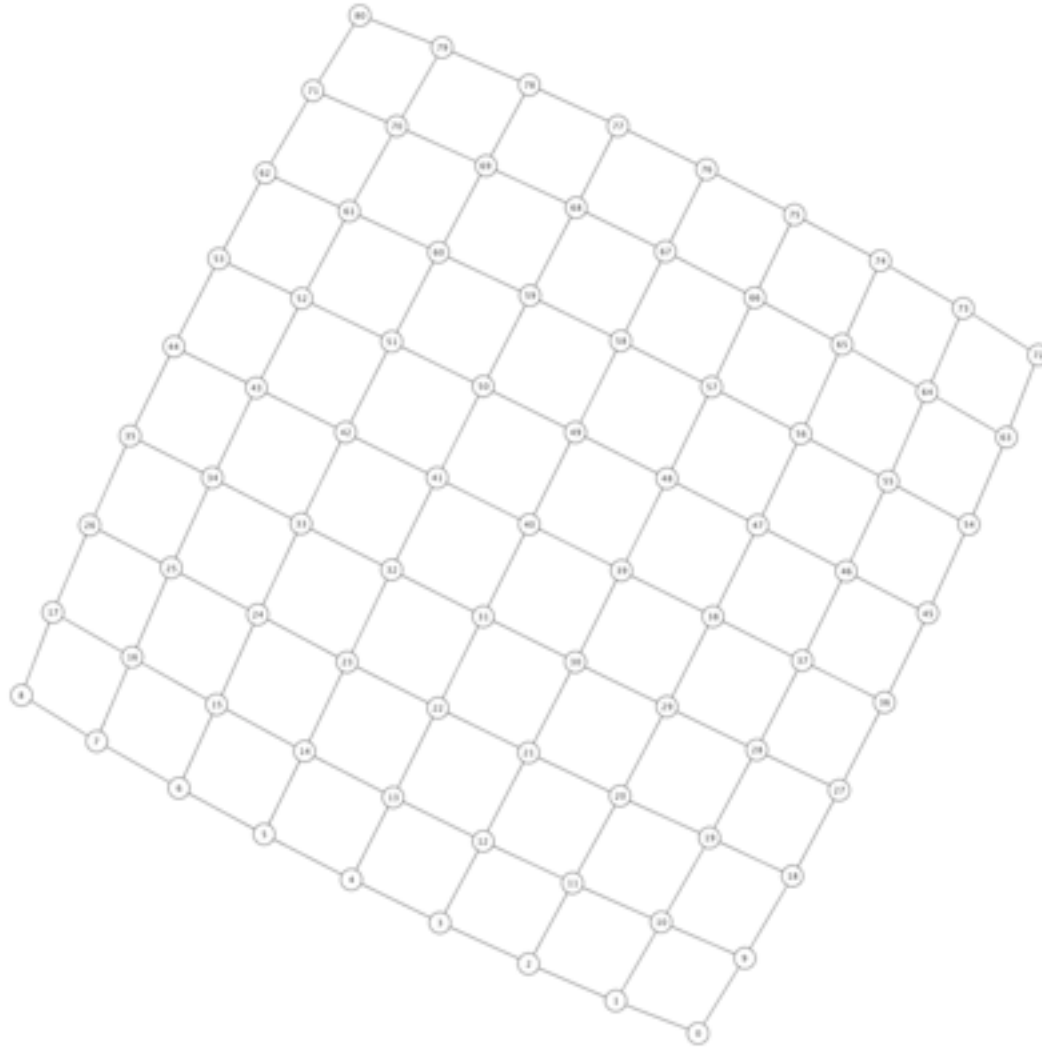
$Z_2$  symmetry (flip all +1 to -1, and vice-versa)

# Erdős-Renyi Random Graph





# Lattice



# Ising Summary

## What is your symmetry?

Each node can be “up” or “down”, and the cost is invariant if you flip everyone at once  
( $Z_2$  symmetry)

## What is your interaction?

pairwise interactions between neighbours on a lattice, that promote sameness  
(nodes want to be in the same state)

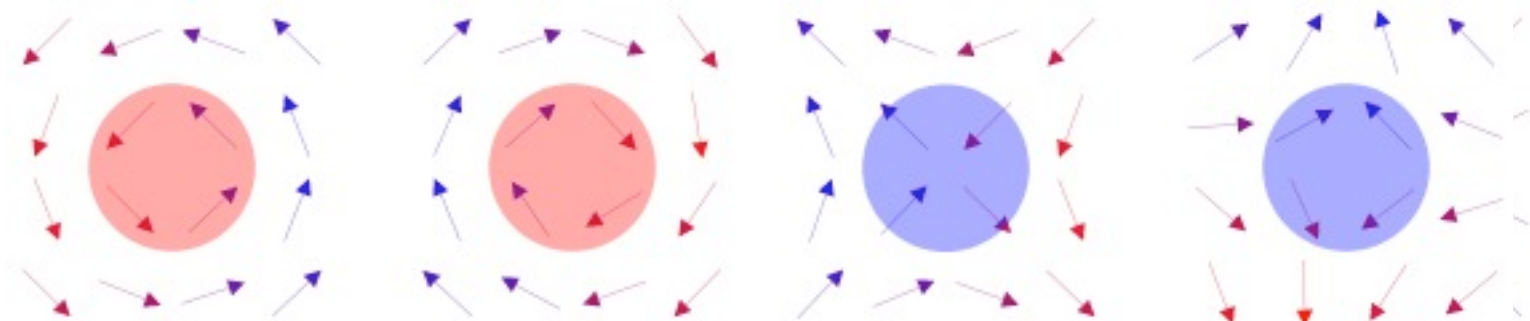
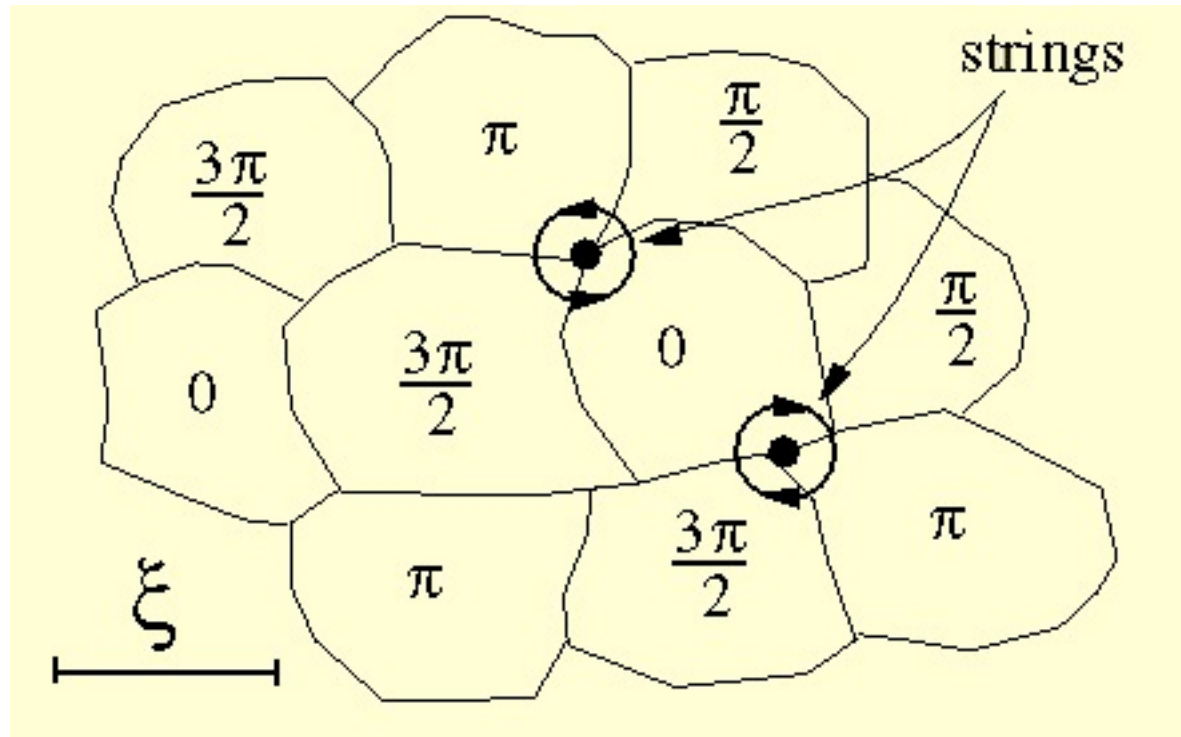
## What is your control parameter?

noise (“temperature”) — higher noise means influence of neighbours  
is small

# [Ising Simulation]



# More Elaborate Symmetries



Vortex Images: Käser, Maier & Rautenkranz (2007)

# Annealing



many defects:  
(quenched)  
strong, but brittle  
— hard to cleave

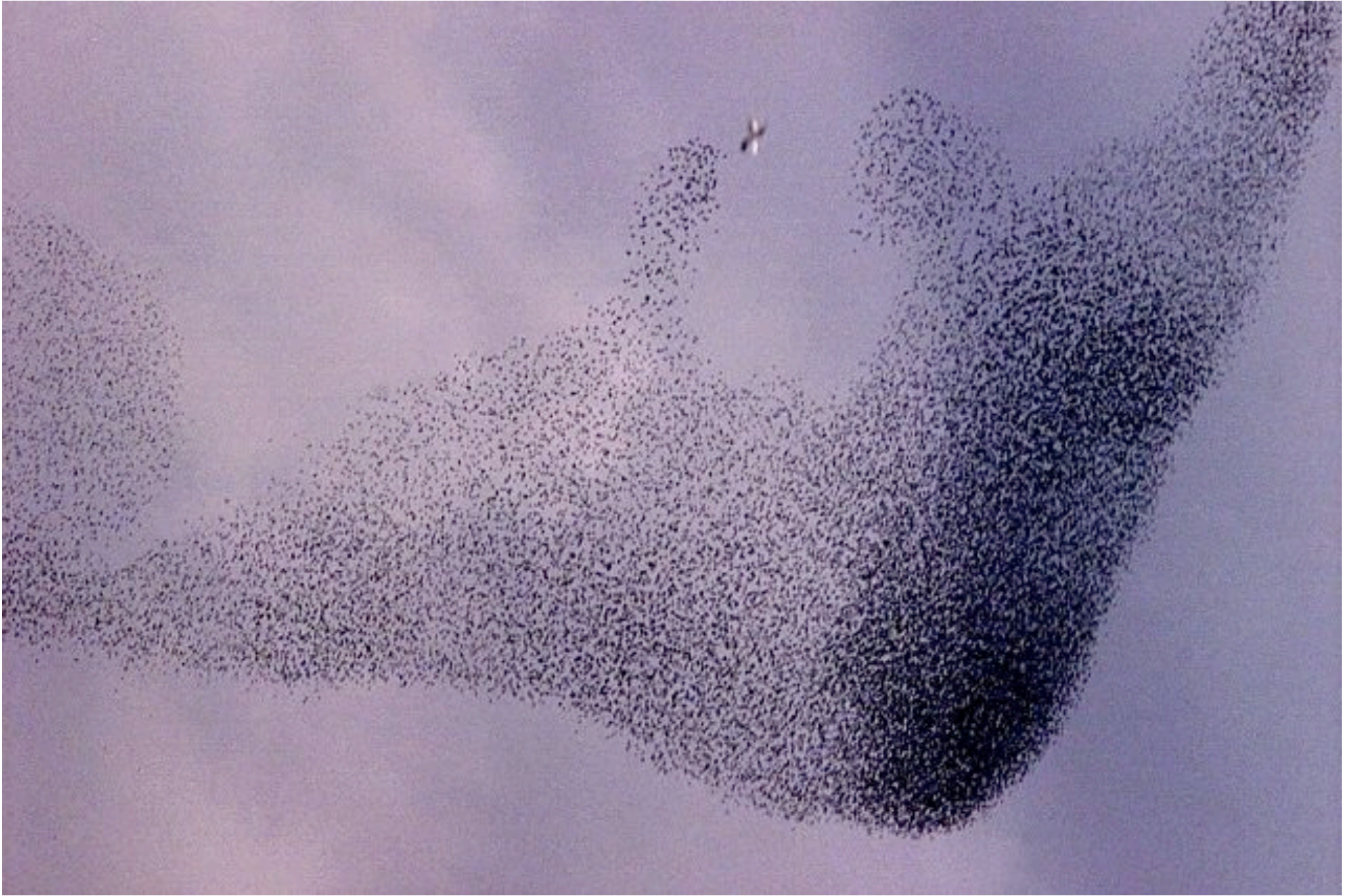
few defects:  
(annealed)  
ductile, flexible



Thursday, June 30, 2011



Thursday, June 30, 2011



Thursday, June 30, 2011





Thursday, June 30, 2011

# What is your symmetry?

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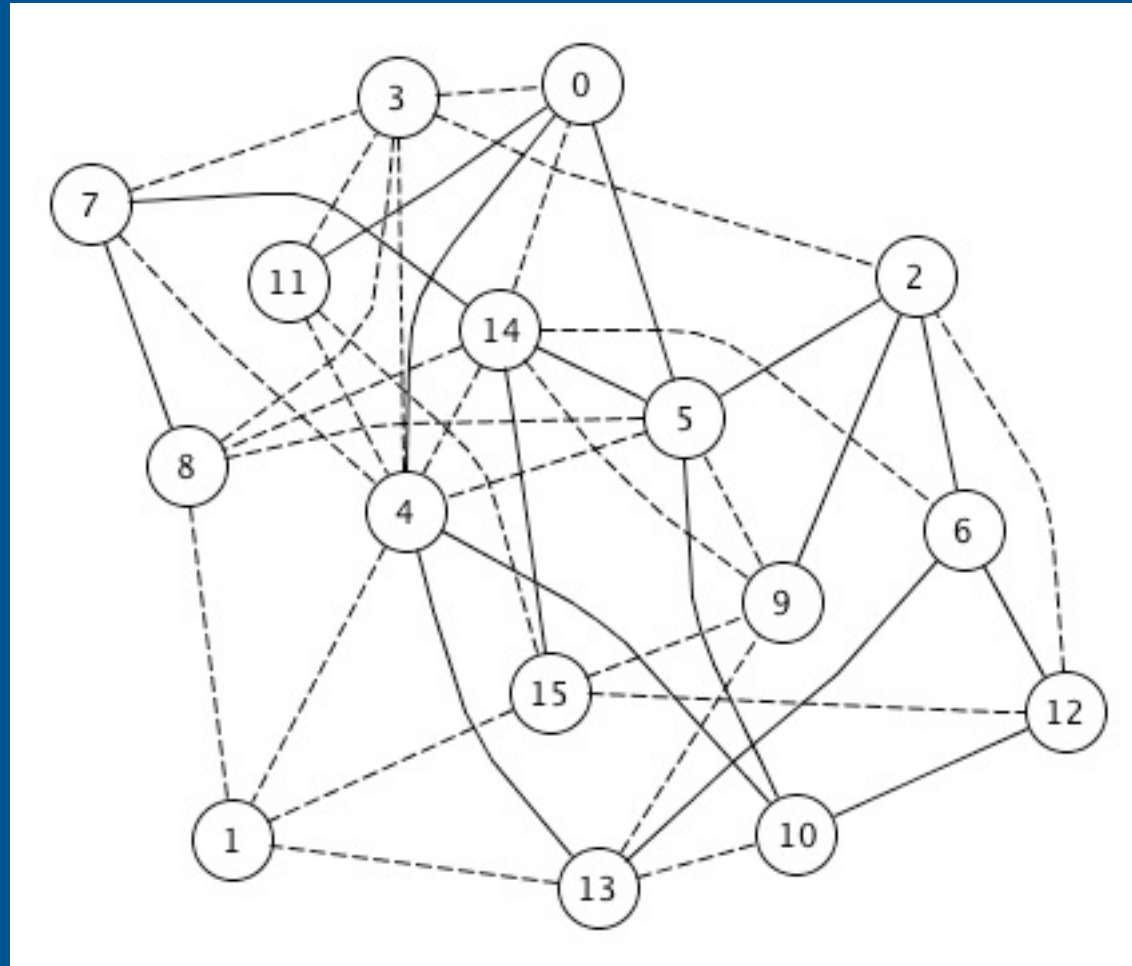
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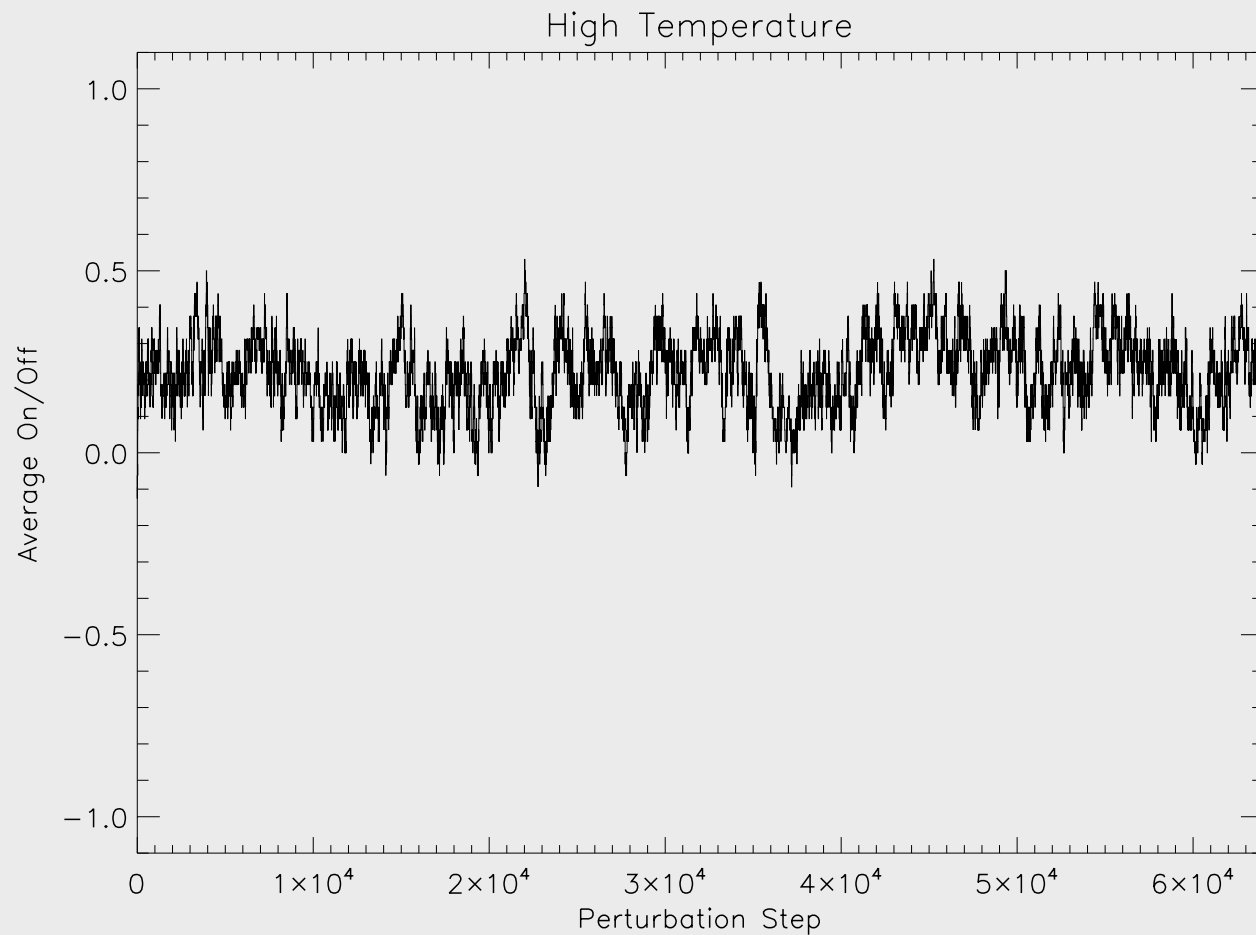
updated worksheet: check wiki, or <http://santafe.edu/~simon/practical.pdf>



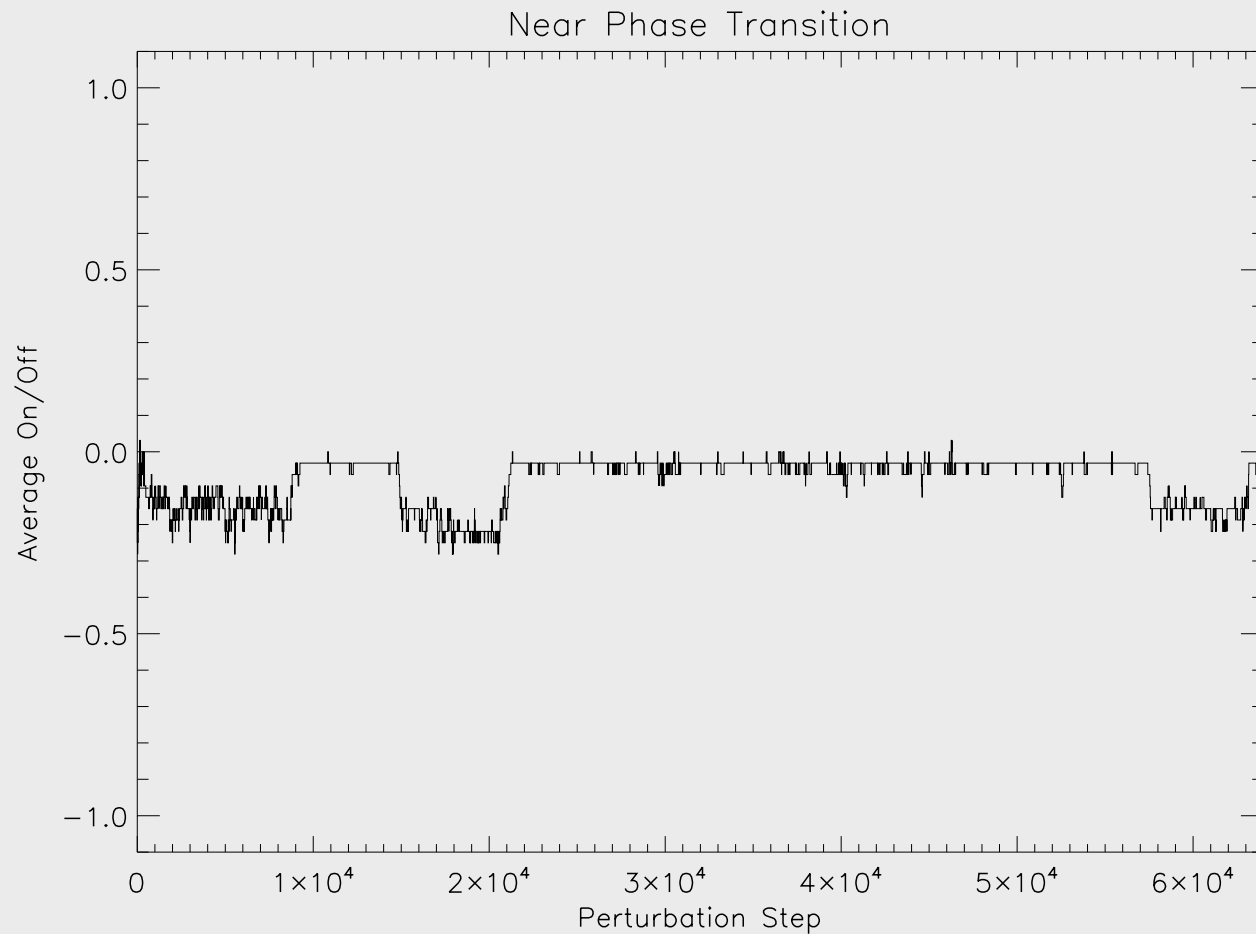
# the Spin Glass



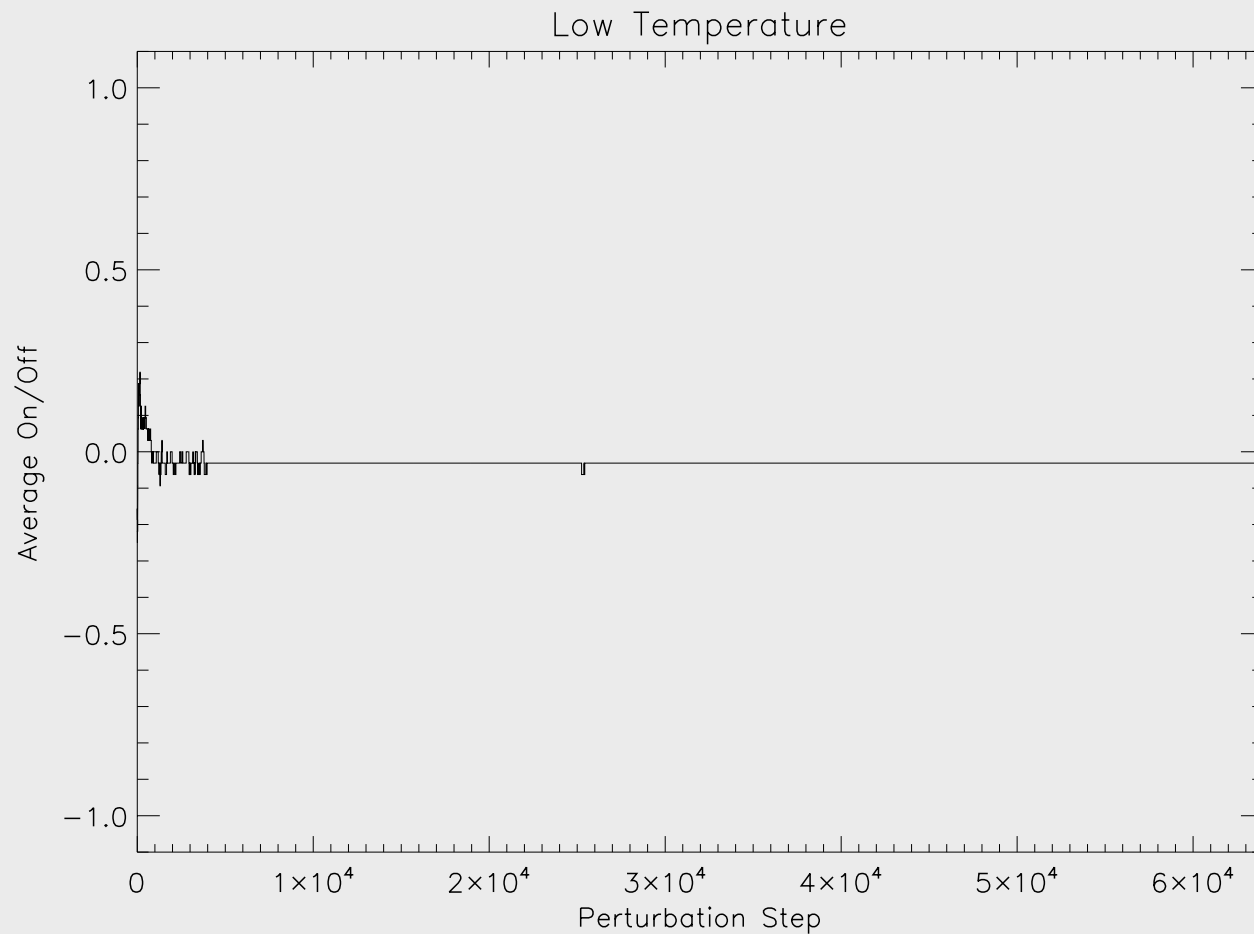
# Memory & the Rugged Landscape



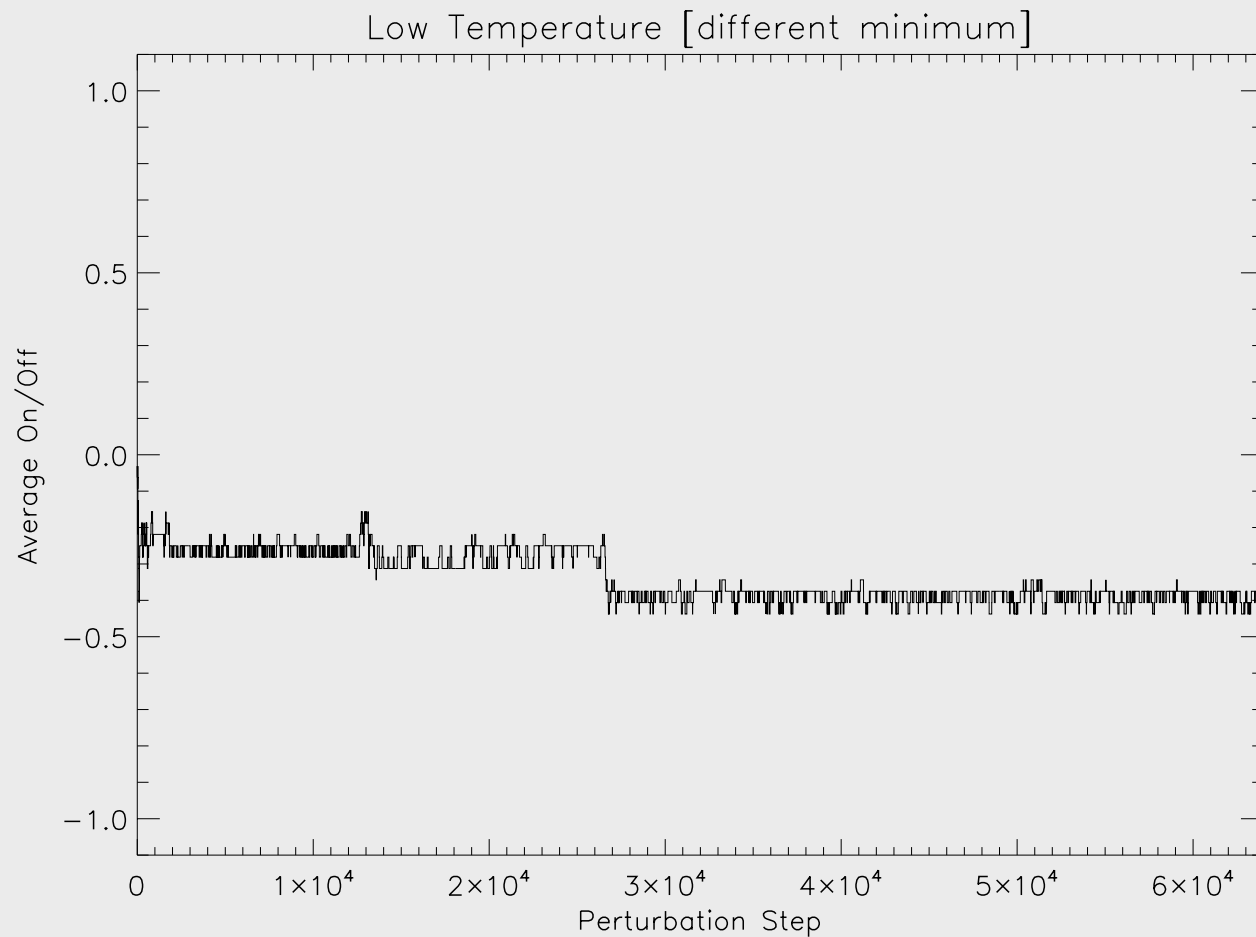
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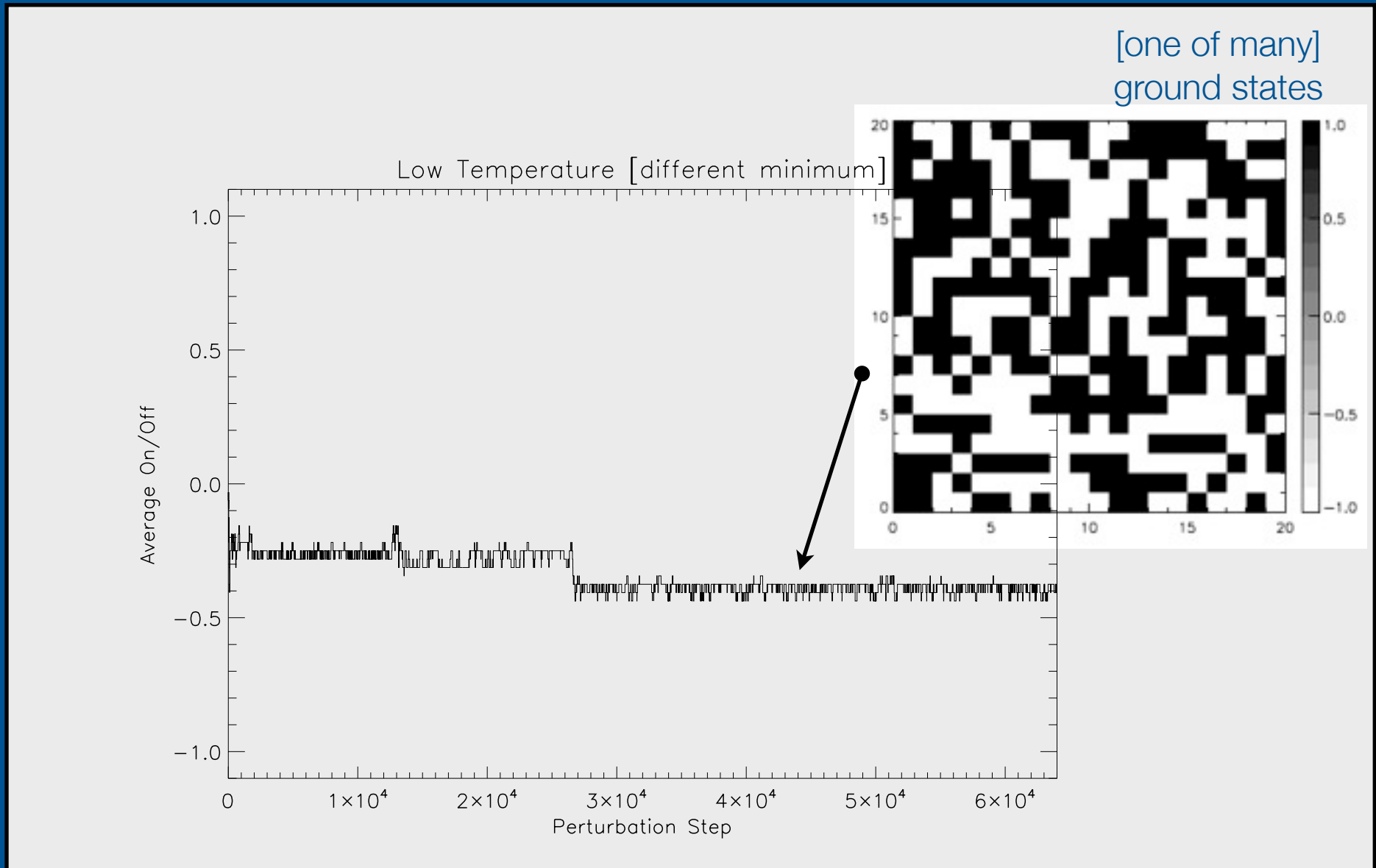
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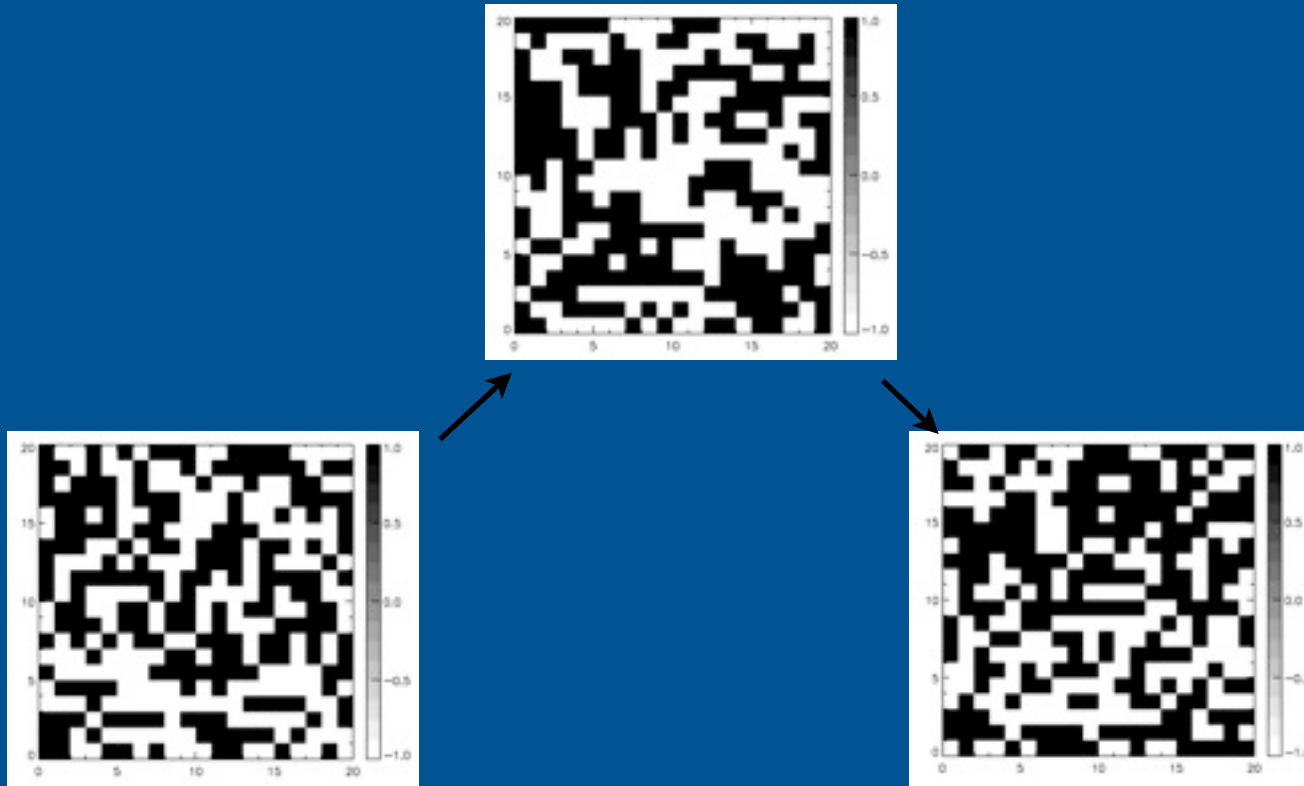


# Memory & the Rugged Landscape



# Sickness & Health

high-temperature



# Salamanders!

## Ising models for networks of real neurons

Gašper Tkačik,<sup>a,c</sup> Elad Schneidman,<sup>a-c</sup> Michael J. Berry II,<sup>b</sup> and William Bialek<sup>a,c</sup>

<sup>a</sup>*Joseph Henry Laboratories of Physics,* <sup>b</sup>*Department of Molecular Biology,*  
and <sup>c</sup>*Lewis-Sigler Institute for Integrative Genomics*

*Princeton University, Princeton, New Jersey 08544 USA*

(Dated: February 4, 2008)

Ising models with pairwise interactions are the least structured, or maximum-entropy, probability distributions that exactly reproduce measured pairwise correlations between spins. Here we use this equivalence to construct Ising models that describe the correlated spiking activity of populations of 40 neurons in the retina, and show that pairwise interactions account for observed higher-order correlations. By first finding a representative ensemble for observed networks we can create synthetic networks of 120 neurons, and find that with increasing size the networks operate closer to a critical point and start exhibiting collective behaviors reminiscent of spin glasses.

PACS numbers: 87.18.Sn, 87.19.Dd, 89.70.+c

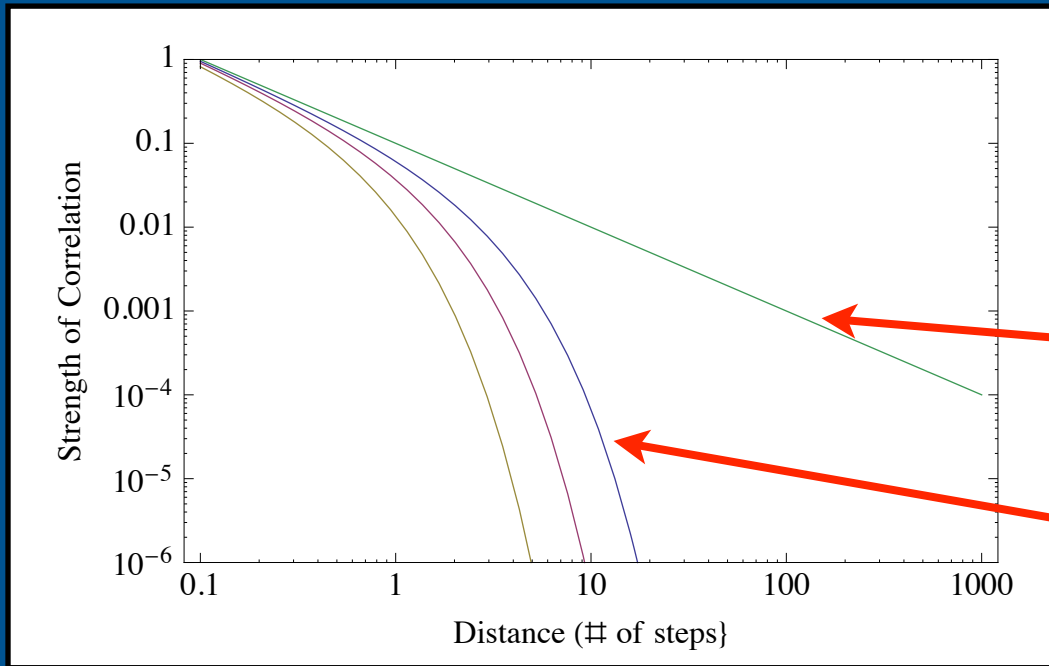
Physicists have long explored analogies between the statistical mechanics of Ising models and the functional dynamics of neural networks [1, 2]. Recently it has been suggested that this analogy can be turned into a precise mapping [3]: In small windows of time, a single neuron  $i$  either does ( $\sigma_i = +1$ ) or does not ( $\sigma_i = -1$ ) generate an action potential or “spike” [4]; if we measure the mean

of Refs [3, 7]. Under these conditions, cells within  $\sim 200 \mu\text{m}$  of each other have  $\sigma_i$  drawn from a homogeneous distribution; the correlation declines at larger distance [8]. This correlated network contains  $N \sim 200$  neurons, of which we record from [9]; experiments typically run for  $\sim 1$  hr [10].

The central problem is to find the magneti



# Critical Opalescence



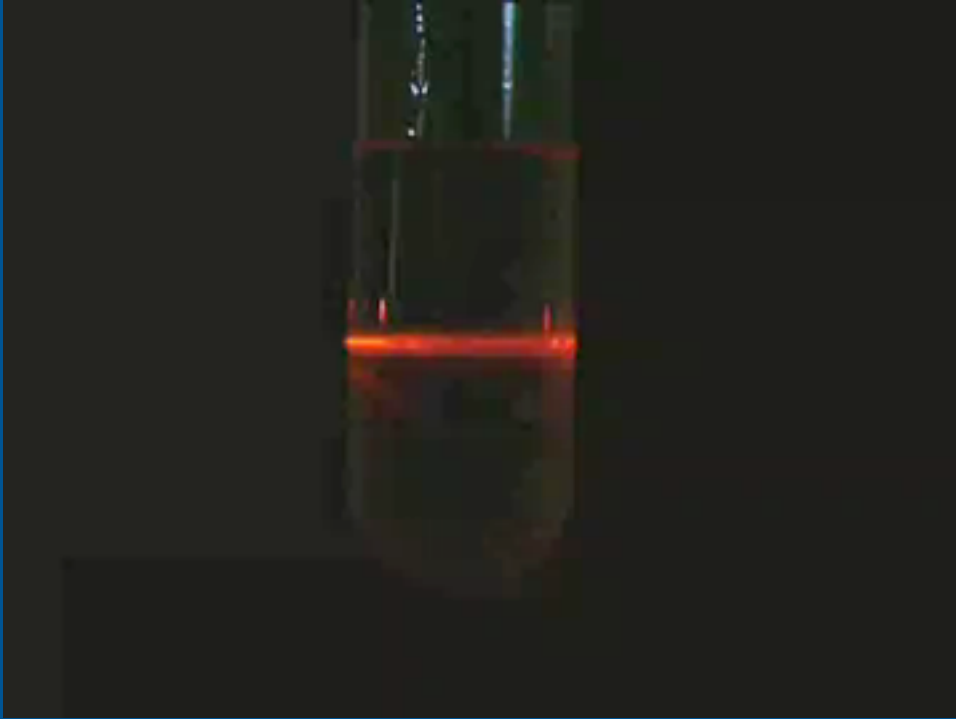
$$e^{n_{\text{steps}}(\beta - \alpha)}$$

Balance! (Power law)

$\alpha, \beta$  not balancing  
above critical temperature

# Critical Opalescence

# Critical Opalescence



# Critical Opalescence

